

$$\nabla \cdot \frac{1}{\rho_e(T)} \nabla V = 0 \quad (6)$$

and current and voltage are related by the vector equivalent of Ohm's Law (Ref. 6)

$$\rho_e J = \frac{\partial V}{\partial r} \quad (7)$$

where J is current density (amp m^{-2}). Solution of Laplace's equation yields the scalar potential field in the material, from which the current density and subsequent Joule heating, q_e , can be derived. The solution requires boundary condition specifications such as $V(r_w, t) = V_w$ (a fixed boundary voltage) and Ohm's law for current flow normal to a boundary

$$J_n = \frac{1}{\rho_e} \frac{\partial V}{\partial n} \Big|_{r_w} \quad (8)$$

where J_n is current density normal to a boundary located at r_w . Finally, note that since Joule heating appears as a source term, the heat equation is coupled with Laplace's equation (which involves temperature-dependent resistivity) and the two must be solved simultaneously, generally by iteration.

It must be noted that the heating phenomena described above are of primary importance in the early (and most time-consuming) phase of wire disintegration. Following melting, magnetic pinch forces rapidly reduce the cross-sectional area of the conductor (Ref. 7), joule heating rapidly increases and expulsion of the melt results. At this point, current conduction is maintained by an arc that grows in length as the wire disintegration process evolves to completion.

There is some prior literature dealing with arc starting in the GMAW process. Melton (Ref. 8) briefly describes the application of the action integral to the analysis of the initial phases of GMAW arc starting and mentions the importance of a sharply pointed tip in rapid arc initiation. Dilthey (Ref. 9) describes a special wire feeder that performs a touch-retract function to reliably start the GMA welding arc. Manz also describes the application of a "sensitivity integral" to the prediction of arc starting times (Ref. 10). Finally, Lenivkin (Ref. 11) reports some experimental current/voltage measurements in the case of arc starting and describes a simple, one-dimensional simulation to represent the initial phases of the arc initiation process.

Experimental Procedure

A Hobart MegaPulse 450 power source/wire feeding system was used to conduct a series of arc starting experiments. The power supply was operated in constant voltage (*i.e.*, nonpulsed) mode. The welding conditions are summarized

in Table 1. Arc starts were made with a matrix of voltage and wire feed speed settings and wire conditions (the end of the wire was either in a freshly clipped state or unclipped and balled from the previous weld). The arc current and voltage were sampled at a 4500-Hz rate using a computer data acquisition system. The arc voltage was measured between the contact tube and workpiece and the arc current was measured with a shunt resistor in series with the welding gun lead. The wire feed speed was monitored by sampling the voltage output of a tachometer in contact with the wire near the motor end of the torch liner. High-speed video images of the arc starts were taken using a Kodak EktraPro 4540 high-speed camera operating at a rate of 4500 frames per second. A 15-W copper vapor laser strobe was used to enhance image quality, allowing clear visualization of the arc start phenomena. The power supply contactor, wire feed motor and welding gun travel motor were all started simultaneously at the beginning of each weld run. Since a delay time was required for the wire to actually contact the base metal, the travel axis was moving at the time of wire contact.

Experimental Results

The experimental matrix of wire feed speed, voltage and wire condition is shown in Table 2. Also included in this table are the actual wire feed speed (measured at the time the wire first contacted the base metal) and the value of the action integral for the initial breakdown of the wire. The voltage and wire feed settings used in these tests correspond to short-circuiting transfer (weld no. 4–17) and spray transfer (weld no. 18–33). The action integral was calculated by Euler integration of the square of the time samples of the weld current. Only that portion of the current waveform corresponding to the first "explosion" of the wire extension was included in the integral. The beginning of the integration interval was identified as the first nonzero current value. The end of the integration interval was identified as the point where the voltage rapidly increased from short circuit value (typically about 10–12 V) to a larger value. The length of the integration interval was typically 20 ms but was sometimes much lower. This integration included the rise time of the current, which typically took about 10 ms to reach a maximum value. This value would be expected to vary with different power supplies, as would other transient characteristics, such as current fall time during arc extinguishments.

Measured current, voltage and wire

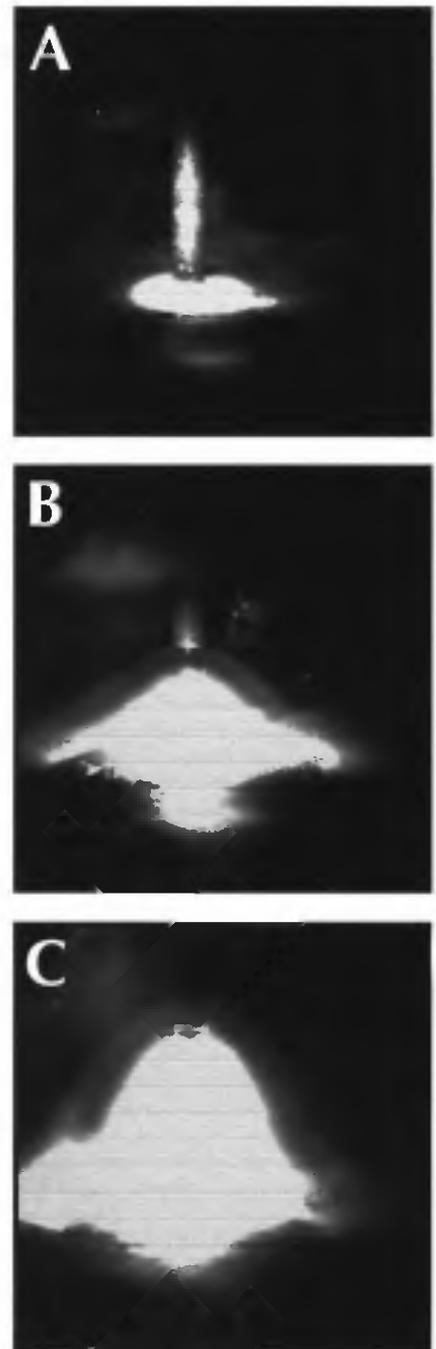


Fig. 2 — Sequence of images corresponding to Weld No. 33.

feed speed waveforms corresponding to Weld No. 33 (where spray transfer settings were used) are shown in Fig. 1. The data indicates a smooth arc start and transition to spray transfer. The current abruptly increased (and voltage decreased) at short circuit and then settled back to the steady state value over a time of about 30 ms. The arc initiation is visible early in the current upslope at a point where the voltage magnitude abruptly increased from near 0 to near 20 V. From Table 2, the action integral of this arc start was relatively low, only about 2.14×10^1 .

high action value starts tended to be sustained, but in some videos, arcing to the contact tube was evident. Presumably, such behavior would promote fusion of the wire to the tube and/or high tube deterioration rates.

Unfortunately, the arc data does not provide clear guidance on how to obtain the smoother starts corresponding to lower action values. Analysis of the data reveals that the lower action integral values were somewhat more likely to be obtained at short circuiting conditions and when the wire was clipped. However, the correlation is not perfect. A number of low action value starts were obtained with no wire clipping and at spray transfer settings. Explanations for this random behavior may lie with inconsistencies in the base metal surface condition (in some cases, arc starts were made at the end of previous passes) or with inconsistency in the shape of the wire end in the unclipped state. In some cases, a large ball was observed on the unclipped wire end while in other cases, the end appeared sharper.

Simulation

To better understand the effects of the shape of the wire tip on arc initiation, physical phenomena occurring in the unfused wire extension during the first phase of arc initiation was simulated. The primary goal of this work was to identify the portion of the wire extension that would first become molten (and hence, fuse and initiate an arc). Evidently, if this occurs at the contact with the base metal, then an arc will initiate at the tip of the wire and a stable arc start may result. If melting first occurs in the middle of the wire extension, then the arc will establish there, which is a less desirable event.

A thermal-electrical simulation of the wire extension was carried out, requiring simultaneous solution of the Laplace and heat equations discussed earlier. Cross sections of the circular geometry and boundary conditions of two simulation cases are depicted in Figs. 6 and 7. In the first case, the wire tip was assumed to be blunt and in good thermal and electrical contact with the base metal. Hence, the lower tip of the wire was assumed to remain at room temperature and at ground potential (0 V) during the arc initiation event. The upper end of the wire extension was assumed to be in good thermal and electrical contact with the contact tube. Thus, the temperature of the upper end of the extension was assumed to remain at room temperature and a total current flow of 1000 A was uniformly distributed over the cylindrical wire cross-section. The second case was identical

Table 3 — Thermophysical Parameters of Steel for Simulation

Heat Capacity ($\text{J kg}^{-1} \text{C}^{-1}$)	750
Thermal Conductivity ($\text{J s}^{-1} \text{m}^{-1} \text{K}^{-1}$)	50
Density (kg m^{-3})	7900
Electrical Resistivity (ohm m)	$15 \times 10^{-8} + 0.004$ $\times 10^{-8} \cdot T (\text{C})$

except the wire was assumed to have a sharp tip, as depicted. Thermo-physical parameters used in the simulation (corresponding to low-carbon steel) are given in Table 3. The primary simulation output was the temperature distribution in the wire at a specified time after the start of current flow (assumed constant at 1000 A). The simulation was performed using a commercially available finite element code (Ref. 12)

A comparison of the simulated temperature fields for two of the cases is shown in Figs. 8 and 9. In both cases, only the tip of the wire is shown so that more details can be resolved. The simulation results show that the temperature distribution in the two wire segments was dramatically different. In the case of the tapered wire, the maximum temperature is achieved in a region very near the tip. At the same point in time, the temperature of the upper section of the wire was still very near to room temperature. This indicates that disintegration of a wire having this shape would first occur near the base metal. On the other hand, the blunt wire achieved melting temperature in a very broad region near the middle of the extension length. Due to effective cooling by conduction to the base metal, the lower end of the wire remained near room temperature. This indicates that a wire having this geometry would first disintegrate over a large region near the center of the extension.

The time taken for the maximum temperature in the simulated wire to reach the melting temperature is also near that observed experimentally. For the cases where the entire wire extension disintegrated, the time taken was typically about 20 ms from the time current first began to flow. In the simulation, the time for the extension to reach melting was about 10 ms. The longer experimental value is explained by the fact that the high short-circuit current levels were not reached instantly; the maximum current levels were typically achieved in about 10 ms — Fig. 3. Also, the wire presumably would disintegrate some time after the instant in which melting is first reached. The sharply tapered wire melted after about 0.1 ms in the simulation. Experimentally, the time taken for

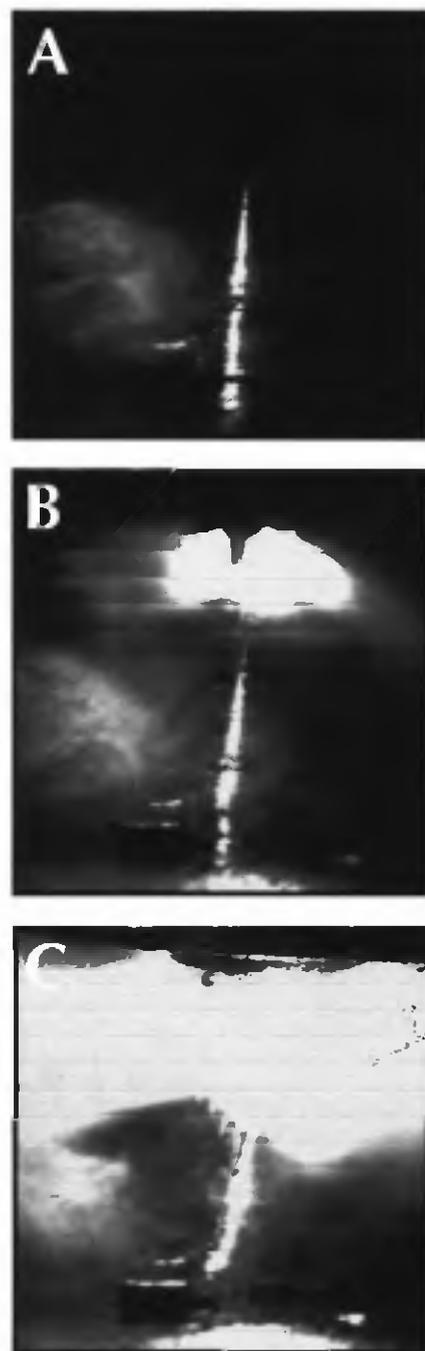


Fig. 4 — Sequence of images corresponding to Weld No. 36.

the wire to fuse was about 2 ms, somewhat longer than the simulation value.

In some cases, the predicted and observed disintegration modes were well matched, but there was "scatter" in this correlation. A number of blunt wires (those having a ball on the end from the previous weld) were observed to disintegrate first at midextension and a number of freshly clipped wires were observed to disintegrate first at the tip. However, quite a number of freshly clipped wires were also observed to disintegrate at midsection and a few unclipped wires

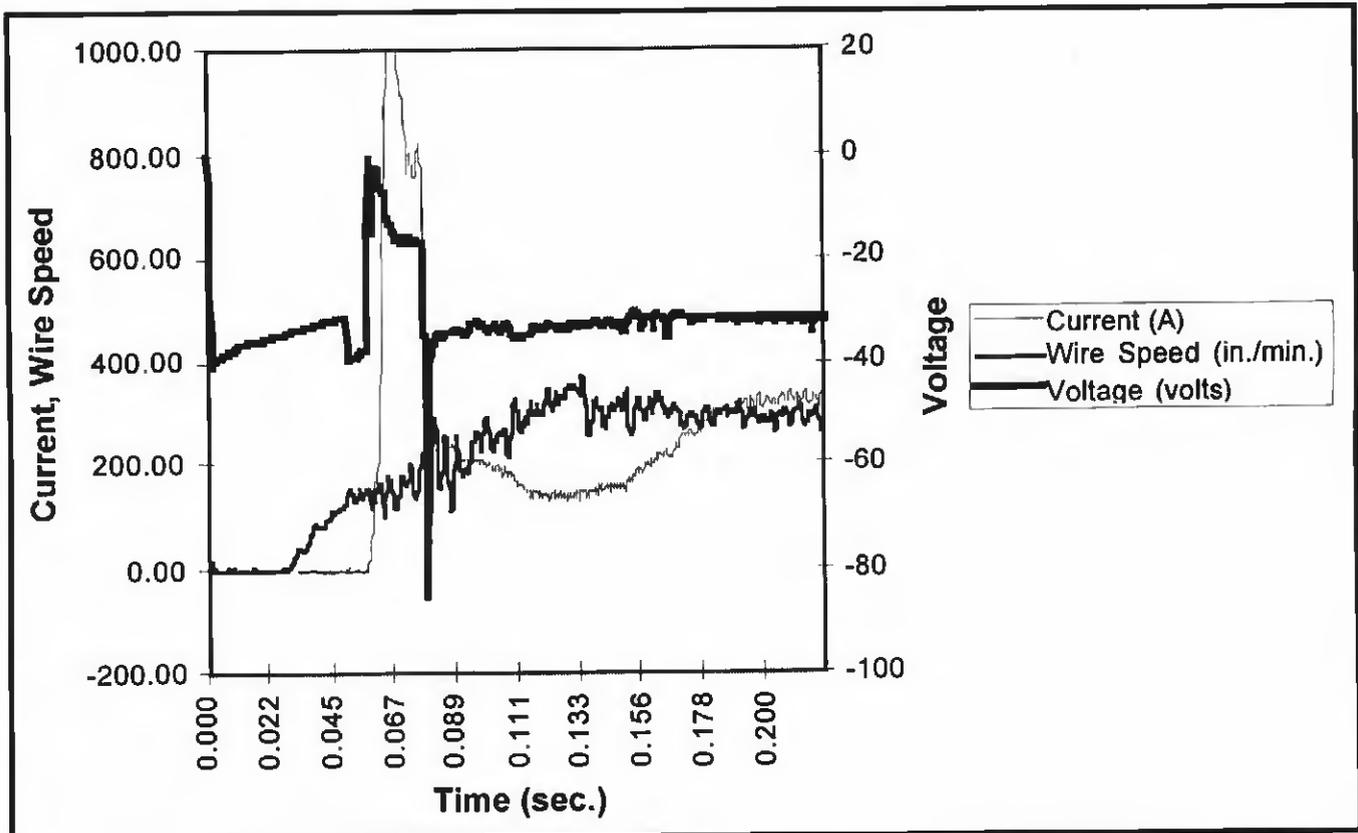


Fig. 5 — Plot of Weld No. 18 current, wire feed speed and voltage waveforms for a case where a violent arc start and subsequent spray transfer were observed.

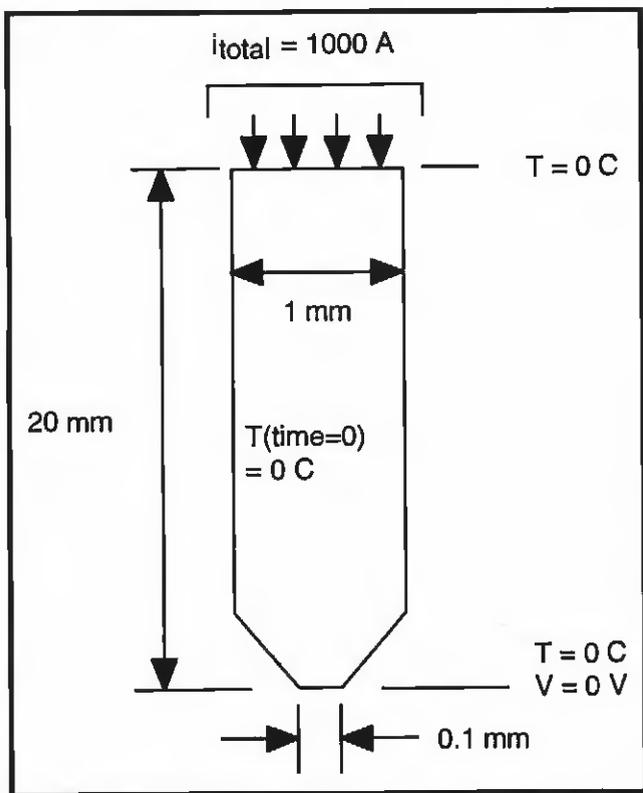


Fig. 6 — Geometry, boundary and initial conditions for tapered wire simulation.

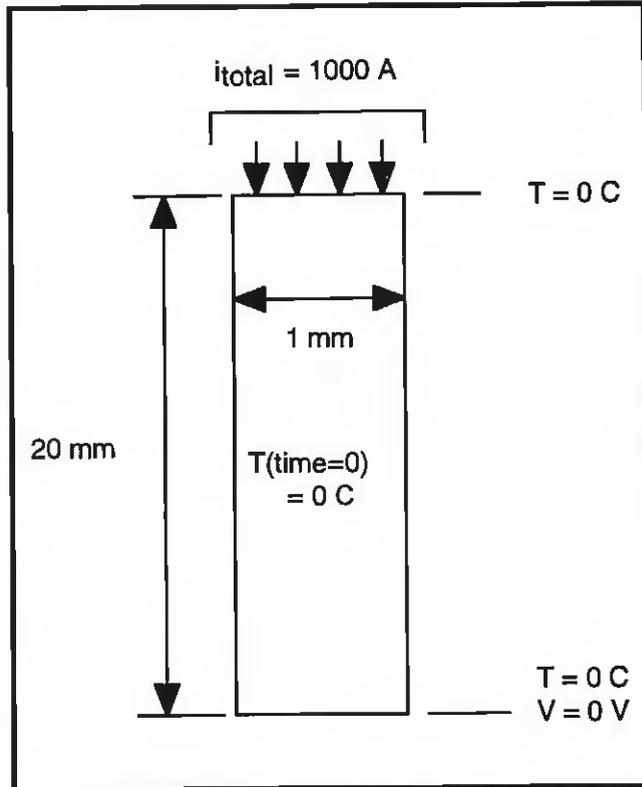


Fig. 7 — Geometry, boundary and initial conditions for blunt wire simulation.

